A Definitive Investigation of the Gas-Phase Two-Center Three-Electron Bond in $[H_2S::SH_2]^+$, $[Me_2S::SMe_2]^+$, and $[Et_2S::SEt_2]^+$: Theory and Experiment

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Abstract: The association products of reactions 1 and 2, examples of two-center three-electron (2c-3e) S : S bonds, were studied in the gas phase. $Me_2S^+ + SMe_2 + M \rightleftharpoons [Me_2S::SMe_2]^+ + M$ (1) and $Et_2S^+ + SEt_2 + M \rightleftharpoons [Et_2S::SEt_2]^+ + M$ (2). The binding enthalpies were determined by measuring the temperature dependence of the equilibrium constants for eqs 1 and 2. The experimental bond enthalpy and entropy of association were determined for reaction 1 at 576 K ($\Delta H^{\circ}_{\text{bond},576} = 111 \pm 2 \text{ kJ/mol}$, $\Delta S^{\circ}_{\text{rxn},576} = -112 \pm 3 \text{ J/mol K}$) and for reaction 2 at 520 K ($\Delta H^{\circ}_{\text{bond},520} = 116 \pm 3 \text{ kJ/mol}$, $\Delta S^{\circ}_{\text{rxn},520} = -132 \pm 5 \text{ J/mol K}$). The calculated bond enthalpies with zero-point and heat capacity corrections are 122 kJ/mol at 576 K at the [PMP4/6-31+G(2df,p)]//MP2/6-31G(d) level and 112 kJ/mol for [Et_2S::SEt_2]^+ at 520 K at the PMP2/6-31G(d)//HF/6-31G(d) level. Three conformations of Et_2S and Et_2S^+ are predicted to be within 4 kJ/mol of each other. The preferred [Et_2S::SEt_2]^+ 2c-3e complex (C_2 symmetry) is formed between C_1 conformers of Et_2S and Et_2S^+. The calculated bond energy for [H_2S::SH_2]^+ is 119.6 kJ/mol using G2 theory. All three 2c-3e complexes were further studied by calculating the lowest optical transition and the hydrogen hyperfine coupling constants.

I. Introduction

The sulfur-sulfur bonds in $[H_2S:SH_2]^+$, $[Me_2S:SMe_2]^+$ ((DMS)₂⁺), and $[Et_2S:SEt_2]^+$ ((DES)₂+) are examples of twocenter three-electron (2c-3e) bonds^{1,2} which were first described in 1931 by Linus Pauling.³ A simple molecular orbital representation of this form of bonding is shown by the potential energy diagram.



Despite the interest in S: S 2c-3e bonds, 4^{-23} very few experiments have directly probed the nature of these bonds. The

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complexes are easily formed in solution and give rise to distinctive ESR^{24} and visible²⁵ spectra from which indirect information on the nature of the 2c-3e S.: S bond can be obtained.²⁶ On the other hand, 2c-3e bonds have been the subject of numerous theoretical investigations.^{27–32} Our interest is to directly determine the gas-phase bond energy. Our earlier efforts³³ were hampered by a small temperature range in the mass spectral study. It was possible to place limits on the bond energy, but a definitive measurement was not possible.

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Figure 1. Schematic cross-section of the new ion source showing internal drift guard rings. The electron entrance plate is electrically insulated from the source block and heat shield by mica gaskets. The drift length is 2.0 cm, while the inside drift guard diameter is 3.0 cm.

In the present paper, we present a more thorough study of the measured and theoretically predicted sulfur-sulfur interactions in $[Me_2S \therefore SMe_2]^+$ and $[Et_2S \therefore SEt_2]^+$. $[H_2S \therefore SH_2]^+$ was only theoretically studied, since the proton-bound dimer rather than the 2c-3e complex is formed in the gas phase at high ion source pressures. The experimental goals of our studies are to determine ΔG°_{rxn} from equilibrium studies and the bond enthalpy (ΔH°_{bond}) from the temperature dependence. Theoretically, the goals include calculating the bond energy, entropy of association, geometry, electron densities, λ_{max} , and hyperfine coupling constants.

The previously reported experimental estimate of the bond enthalpy for $[Me_2S \therefore SMe_2]^+$ was derived from high-pressure mass spectrometric studies of the equilibrium constant for reaction 1 at a single temperature (hence a single determination of ΔH°_{rxn} combined with an estimate of the entropy of reaction).³³ Due to the relatively strong S. S bond, high temperatures were required for the equilibrium study, and we were only able to make measurements at the high-temperature limit of our apparatus. Thus, the experimental bond enthalpy estimate for the 2c-3e interaction in $[Me_2S \therefore SMe_2]^+$ was in the range 100–111 kJ/mol. Ab initio calculations at the [PMP2/ 6-31G(d)]//3-21G(d) level gave a bond energy of 109.6 kJ/mol³³ (the brackets indicate that the additivity approximation³⁴ was used).

In earlier experiments, the enthalpy of reaction for the iodine iodine interaction in $[CH_3I.:ICH_3]^+$ was estimated to be in the range -96 to -109 kJ/mol.^{35,36} This range represents about 65% of the iodine—iodine interaction in I_2^{37} and is in excellent agreement with the result reported for the iodine—iodine 2c-3e bonding interaction of 102 kJ/mol in $I_2^{-.27}$ We believe that these studies on $[Me_2S::SMe_2]^+$ and $[CH_3I::ICH_3]^+$ represented the first gas-phase experimental binding results for wellcharacterized 2c-3e bonds between heteroatoms in organic molecules. The present study offers more detailed experimental bond information as well as new theoretical insights into this type of bonding.

II. Experimental Section

A redesigned ion source (Figure 1) capable of a much larger temperature range (90-650 K) and improved thermalization of all

source parts was used in the present studies. The source has been described elsewhere in detail.³⁸ Of importance here is the fact that the ion exit aperture is mounted **inside** a large copper block. This ensures that the entire region near the ion exit is at the same temperature as the source block itself. The electron filament was shielded first by a biasing plate and second by a thermal shield, thus reducing the "hot spot" near the ionizing filament. As in previous designs, internal drift guard rings shaped and controlled the electric fields (Figure 1).^{38–40}

It should be noted that the ionization voltage was very low (9-15 V) which means the electrons do not penetrate far into the ion source but rather form ions only near the electron entrance hole. All of the charges were pulled through the entire length of the ion source by the extraction voltage and underwent many thousands of collisions during the transit time of >100 μ s.⁴¹ Hence, the ions were well thermalized with little initial excess energy, resulting in reduced fragmentation and therefore simple high-pressure mass spectra.³⁸⁻⁴²

Our arrangement may be contrasted to high-pressure ion sources that use conventional 90° electron entrance—ion exit geometry. At voltages typical for these geometries (>1000 V), electrons produce an electrical plasma and are scattered throughout the source, including the region near the ion exit slit, where they cause ionization.⁴³ Residence time distributions (RTD's) for these sources, therefore, usually show reactant ion signals near $t = 0.^{44}$ These ions have not undergone many switching reactions and are not in equilibrium. Hence, RTD's from these sources do not achieve constant product-to-reactant ion ratios (and equilibria) until some time after the ionizing pulse.^{43–45}

An earlier published study⁴⁶ on the bonding in $H_3O^+CO_2$ and $H_3O^+N_2O$ showed that our new source gave results that compare very well with results from other groups using different ion source configurations. In the case of $H_3O^+CO_2$, other groups obtained bond enthalpies that ranged from 59.8 to 64.0 kJ/mol and entropies of association that ranged from -86.6 to -102.9 J/mol K.⁴⁷ With our source we obtained an experimental bond enthalpy of 56.1 kJ/mol and an entropy of association of -82.8 J/mol K, which agrees very well with a high level ab initio study.⁴⁸ In the case of $H_3O^+N_2O$, the result from unpublished work by Szulejko and McMahon was 66.1 kJ/mol for the experimental bond enthalpy and -93.7 J/mol K for the entropy of association.⁴⁷⁴ Our experimental bond enthalpy was 65.3 kJ/mol, and the entropy of association was -100.8 J/mol K,⁴⁶ which again agrees very well with a published ab initio study.⁴⁹

In the present studies on $[Me_2S \therefore SMe_2]^+$, total ion source pressures ranged from 0.50 to 0.73 Torr. The Me₂S (DMS) concentration covered the range 5.9–14.8 mol% in CF₄ bath gas. The $[Et_2S \therefore SEt_2]^+$ experiments were carried out in CF₄ bath gas with total ion source pressures ranging from 0.31 to 0.45 Torr and the Et₂S (DES) concentration ranging from 1.9 to 4.7 mol%. The reactant gas, Me₂S

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Table 1. Total Energies (hartrees) and Zero-Point Energies (kJ/mol) for Various Species^a

				HF/6-31G(d) ^b			1	/MP2/6-31G(d)		
	sym	sta	HF/a	PMP2/a	ZPE ^c /a	MP2/a	PMP2/a	PMP4/a	PMP2/b	$G2^d$
SH ₂ 1	C_{2v}	¹ A ₁	-398.66732	-398.78821	43.2 (0)	-398.78841	-398.78841	-398.81200	-398.85551	-398.93072
SH ₂ ⁺ 1 ⁺	C_{2v}	${}^{2}B_{1}$	-398.32699	-398.42719	42.6 (0)	-398.42562	-398.42751	-398.45100	-398.48129	-398.54744
$(SH_2)_2^+$	$C_{2\mathrm{h}}$	${}^{2}\mathbf{B}_{u}$	-797.02610	-797.26395	100.2 (0)	-797.26162	-797.26461	-797.31048	-797.38999	-797.52371
$H_2S \cdot 2H_2O$	C_{s}	¹ A′	-550.69777°	-551.19153°	172.9 (0) ^e	-551.20092	-551.20092	-551.24460	-551.42976	
$H_2S \cdot 2H_2O^+$	$C_{\rm s}$	² A″	-550.41162	-550.89446	178.3 (1)	-550.89288	-550.89468	-550.93621	-551.10483	
$H_4S_2 \cdot 4H_2O^+$	C_{i}	${}^{2}A_{u}$	-1101.14725	-1102.13918	367.1 (0)	-1102.14204	-1102.14488	-1102.22892	-1102.59006	
DMS 2	C_{2v}	$^{1}A_{1}$	-476.73533	-477.12066	214.1 (0)	-477.12111	-477.12111	-477.17503	-477.26052	
DMS ⁺ 2 ⁺	C_{2v}	${}^{2}B_{1}$	-476.45471	-476.81704	211.6 (0)	-476.81573	-476.81776	-476.87232	-476.94651	
$(DMS)_2^+$	$C_{2\mathrm{h}}$	${}^{2}\mathbf{B}_{u}$	-953.21557	-953.98535	433.1 (0)	-953.98357	-953.98657	-954.09277	-954.26001	
DES 3a	C_{2v}	$^{1}A_{1}$	-554.80562	-555.45409	375.4 (0)					
DES 3b	C_2	ιA	-554.80429	-555.45400	375.4 (0)					
DES 3c	$C_{\rm s}$	¹A′	-554.80253	-555.45201	374.9(1)					
DES 3d	C_1	۱A	-554.80496	-555.45402	375.4 (0)					
DES ⁺ 3a ⁺	C_{2v}	${}^{2}B_{1}$	-554.53569	-555.15955	372.6 (0)					
DES ⁺ 3b ⁺	C_2	^{2}B	-554.53390	-555.15908	375.4 (0)					
DES ⁺ 3c ⁺	$C_{\rm s}$	$^{2}A'$	-554.53247	-555.15731	374.1 (1)					
DES ⁺ 3d ⁺	C_1	^{2}A	-554.53475	-555.15928	373.5 (0)					
$(DES)_2^+$	C_{2h}	${}^{2}\mathbf{B}_{u}$	-1109.35868	-1110.65790	754.8 (3)					
$(DES)_2^+$	C_2^{-1}	${}^{2}\mathbf{B}^{-}$	-1109.36159	-1110.66046	754.6 (0)					

^{*a*} Basis set a is 6-31G(d). Basis set b is 6-31+G(2df,p). ^{*b*} Spin-squared value for open-shell species were between 0.76 and 0.78. ^{*c*} Zero-point energy (kJ/mol) with number of imaginary frequencies in parentheses. ^{*d*} G2 method, ref 54. ^{*c*} The HF/6-31G(d) geometry is assumed to be SH₂ plus (H₂O)₂.

or Et₂S, was mixed with the bath gas in the ion source inlet system by introducing the gases through separate leak valves. Gas pressures were measured with an MKS capacitance manometer. Mass spectra were recorded using conventional analog recording techniques, while residence time distributions and ion intensity measurements for the determination of equilibrium constants were collected using previously described pulse-counting techniques.^{39,40}

Computational Methods. After comparison of MNDO, ^{50a} AM1, ^{50b} and PM350c with ab initio results on the thiirane 2c-3e complex in a previous publication,⁵¹ MNDO was chosen in the present study to survey the Et_2S and $[Et_2S::SEt_2]^+$ stationary points. The MNDO geometries were used as starting positions of ab initio optimizations using the Gaussian92 program system.⁵² All geometries were optimized at the HF/6-31G(d) level, and frequencies were calculated at that level to confirm the nature of the stationary point and to make corrections for zero-point energies (0.9 weighting factor). Single-point calculations were made at the PMP2/6-31G(d) level where "P" indicates that the effect of spin contamination has been projected out of the MP energy.53 Additional optimizations at the MP2/6-31G(d) level were carried out for all species except Et₂S, Et₂S⁺, and [Et₂S∴SEt₂]⁺. At MP2/6-31G-(d) geometries, single-point calculations were made at the PMP4/6-31G(d) and PMP2/6-31+G(2df,p) levels and combined using the additivity approximation³⁴ to estimate relative energies at the [PMP4/ 6-31+G(2df,p)] level. For H_2S , H_2S^+ , and $[H_2S \therefore SH_2]^+$, G2 theory, which estimates energies at the [QCISD(T)/6-311+G(3df,2p)] level with corrections for zero-point energy and higher order effects, was applied. G2 theory⁵⁴ has been shown to yield results of chemical accuracy and has been applied to several studies of sulfur-containing species.55-57 The lowest excitation energy in the 2c-3e bonded species was calculated

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using CI singles method at the CIS/6-31G(d) level.⁵⁸ Thermodynamic corrections (to determine standard enthalpy, entropy, and free energy) were made using uncorrected vibrational frequencies at the HF/6-31G-(d) level with standard methods⁵⁹ unless otherwise noted. Absolute and relative energies for all molecules and conformers are presented in Tables 1 and 2, respectively, while calculated geometries are given in Figure 2. Experimental microwave structures for SH₂, SH₂⁺, and DMS are also included in Figure 2.⁶⁰

Thermodynamic Equilibrium Measurements. The thermodynamic bond energies for $[Me_2S :: SMe_2]^+$ and $[Et_2S :: SEt_2]^+$ were determined from equilibrium measurements of the reversible reactions 1 and 2.

$$\operatorname{Me}_2 S^+ + \operatorname{SMe}_2 + M \rightleftharpoons [\operatorname{Me}_2 S : : \operatorname{SMe}_2]^+ + M$$
 (1)

$$\operatorname{Et}_{2}S + \operatorname{SEt}_{2} + M \rightleftharpoons \left[\operatorname{Et}_{2}S \therefore \operatorname{SEt}_{2}\right]^{+} + M$$
 (2)

Specifically, the van't Hoff equation (eq 3) yields the enthalpy and entropy of reaction.

$$\ln K_{\rm p} = -\Delta H^{\circ}_{\rm rxn}/RT + \Delta S^{\circ}_{\rm rxn}/R \tag{3}$$

For association reactions, the enthalpy of reaction is the negative of the bond enthalpy. The bond energy (D°_{\circ}) is equal to the bond enthalpy at 0 K and can be related to the bond enthalpy at a different temperature by eq 4 where ΔC_p is the difference in heat capacities between reactants and products.

$$D^{\circ}_{o} = -\Delta H^{\circ}_{T,rxn} + \int_{o}^{T} \Delta C_{P}(T) dT \qquad (4)$$

For ion-molecule association reactions, the heat capacity contribution of the cluster is usually larger due to a population of low-lying modes formed between the two reactant moieties; hence, the value for the bond energy is usually larger by about 4-8 kJ/mol than the value for the bond enthalpy. This difference is expected to be smaller the stronger the bond between the two moieties. Statistical thermodynamic heat capacity corrections for systems involving gas-phase ions are often

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Figure 2. Calculated geometries for monomers and 2c-3e complexes at the HF/6-31G(d) level with MP2/6-31G(d) values in parentheses. Experimental microwave deterimations for SH_2 , SH_2^+ , and DMS given in brackets are from ref 60.

not possible since the vibrational frequencies and molecular geometries for ions are difficult to obtain. In the present study the parameters predicted by ab initio calculations are used to obtain the true bond energies.

In order to correctly determine experimental equilibrium constants, it must be demonstrated that the system is in true chemical equilibrium. We applied various tests to ensure that equilibrium exists. We measured RTD's for both reactant and product ions which indicate the time that the particular ions spend in the ion source. When a reaction is in chemical equilibrium, RTD's for reactant and product ions are identical because the charges undergo many switching reactions as they traverse the ion source, resulting in an averaging of the drift properties.^{33,34,38,39} The RTD's were normalized by dividing each RTD by the maximum intensity (count). Overlapping RTD's clearly demonstrate that the product-to-reactant ion ratio is constant. Hence, if the RTD's are the same, it can be inferred that the system is at or near equilibrium. Apparent equilibrium constants were also measured as a function of composition, total ion source pressures, and ion extraction fields. The equilibrium data were extrapolated to zero extraction fields.

III. Results and Discussion

 $[H_2S::SH_2]^+$. There have been several previous theoretical calculations on the 2c-3e $[H_2S::SH_2]^+$ complex.^{27–32,61} Recent binding energies range from 111³¹ to 124²⁷ kJ/mol while the S::S distance is about 2.8 Å. Gill et al. have summarized the recent work.³² Our G2 binding energy is 119.6 kJ/mol (Table 2) which is intermediate to the predictions at the PMP2/6-31G-(d)+ZPC and [PMP4/6-31+G(2df,p)]+ZPC levels, (114.5 and 123.5 kJ/mol, respectively, Table 2). The gas-phase S::S

^{(61) (}a) Fernández, P. F.; Ortiz, J. V.; Walters, E. A. J. Chem. Phys. 1986, 84, 1653. (b) Ortiz, J. V. Chem. Phys. Lett. 1987, 134, 366. (c) Kikuchi, O. Bull. Chem. Soc. Jpn. 1981, 54, 917.

Table 2.	Relative	Energies	(kJ/mol) of	Various	Species ^a
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		//HF/6-31G(d)			//MP2/0	5-31G(d)		
	HF/a	PMP2/a	+ZPC	PMP2/a	PMP4/a	PMP2/b	[PMP4/b] ^b	+ZPC ^c	G2
SH ₂ 1	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
$SH_2^+ 1^+$	893.5	9 47.8	947.3	947.5	947.8	982.5	982.8	982.3	1006.3
$1 + 1^+$	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
$(SH_2)_2^+$	-83.4	-127.5	-114.5	-127.8	-124.7	-139.7	-136.6	-123.5	-119.6
$SH_2 \cdot 2H_2O$	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
$SH_2 \cdot 2H_2O^+$	751.3	779.9	784.8	804.0	809.6	853.1	858.7	863.6	
$1(aq) + 1^+(aq)$	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	
$(SH_2)_2^+(aq)$	-99.3	-139.7	-125.4	-129.4	-126.3	-145.6	-142.5	-128.2	
DMS 2	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	
DMS ⁺ 2 ⁺	736.8	797.2	795.0	796.4	794.7	824.4	822.7	820.5	
$2 + 2^+$	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0	
$(DMS)_2^+$	-67.0	-125.1	-118.4	-125.2	-119.3	-139.0	-133.1	-126.4	
DES 3a	0.0	0.0	0.0						
DES 3b	1.9	0.2	0.2						
DES 3c	8.1	5.4	5.0						
DES 3d	1.8	0.2	0.2						
$DES^+ 3a^+$	708.7	773.4	770.9						
DES ⁺ 3b ⁺	713.4	774.5	774.5						
$DES^+ 3c^+$	717.1	779.2	778.0						
DES ⁺ 3d ⁺	711.2	774.0	772.3						
$3a + 3a^{+}$	0.0	0.0	0.0						
$3b + 3b^+$	8.2	1.5	4.0						
$3c + 3c^+$	16.5	11.3	12.2						
$3d + 3d^+$	4.2	0.8	1.6						
$(DES)_2^+(C_{2h})$	-45.6	-116.2	-110.1						
$(DES)_2^+(C_2)$	-53.2	-122.9	-117.0						

^{*a*} Basis set a is 6-31G(d). Basis set b is 6-31+G(2df,p). ^{*b*} Additivity approximation is used. See ref 34. ^{*c*} Zero-point corrections are added to [PMP4/b] relative energies.

distance (2.835 Å) is identical to the value calculated by Gill et al.³² at the same level (MP2/6-31G(d)). Since the only experimental observation of [H₂S:.SH₂]⁺ has been in aqueous solution,²⁵ we also carried out calculations in which water molecules were allowed to interact explicitly with each hydrogen. At the MP2/6-31G(d) level, the oxygen-hydrogen distance is 1.626 Å in $H_2S^+ 2H_2O$ and 1.874 Å in $[H_2S^- SH_2]^+ 4H_2O$ (Figure 2), indicative of relatively strong hydrogen bonds. In contrast, in H₂S·2H₂O, one H₂O is weakly hydrogen bonded to one hydrogen of H_2S (O···H 2.080 Å) while the second H_2O is hydrogen bonded to the first H2O. This weak interaction between H₂S and water is in accord with its known limited solubility in water.⁶² The S∴S bond energy remains almost unchanged on going from the gas phase to solution phase (i.e., $H_2S \cdot 2H_2O + H_2S^+ \cdot 2H_2O \rightarrow [H_2S \therefore SH_2]^+ \cdot 4H_2O)$ which suggests that the solvation energy of H_2S^+ and $[H_2S \therefore SH_2]^+$ are nearly the same. While the binding energy increases only 4.5 kJ/mol on going to the aqueous phase, the effect on the S::S bond distance is more pronounced, shortening 0.094 Å. As will be discussed below, this difference will affect its visible spectrum. Schlever and co-workers⁶³ found a similar shortening in the complex H_3BNH_3 ($\Delta BN = 0.05$ Å at MP2/6-31G(d)) when they explicitly optimized the adduct with three waters around the nitrogen end. They suggested that the bond shortening was due to the dielectric reaction field of water (which would increase the dipole moment of H₃BNH₃). In the case of $[H_2S: SH_2]^+ 4H_2O$, the explanation may be due to decreased electrostatic repulsion from a redistribution of positive charge to the solvent. At the UHF/6-31+G(2df,p) level, 0.14e⁻ of charge has been transferred from the four solvent molecules to the 2c-3e cation.64

Olson, L.; Reichel, F.; Kraka, E. Isr. J. Chem. 1993, 33, 369.



Figure 3. Normalized residence time distributions for Me_2S^+ (open diamonds) and $[Me_2S \therefore SMe_2]^+$ (filled circles) which indicate that the reactants and products in reaction 1 are at equilibrium.

 $[Me_2S::SMe_2]^+$. In an earlier paper the ion-molecule reactions responsible for the major peaks in the mass spectrum were identified.³³ Figure 3 shows that RTD's for Me₂S⁺ and $[Me_2S::SMe_2]^+$ collected with low drift potentials are identical. Measured apparent equilibrium constants, K_{app} , must also be invariant to the chemical composition (partial pressures). Experiments at various compositions and total ion source pressures were extrapolated to zero ion extraction volts at each temperature studied. At the lowest Me₂S⁺ concentration, plots of K_{app} vs extraction voltage show a voltage dependence. At these low concentrations, equilibrium is not fully established, and the plots demonstrate kinetic approach to equilibrium as

⁽⁶²⁾ H₂S solubility: CRC Handbook of Chemistry and Physics, 71st ed.; Lide, D. R., Ed.; CRC Press: Boca Raton, FL, 1990–1991; pp 4–68. (63) (a) Bühl, M.; Steinke, T.; Schleyer, P. v. R.; Boese, R. Angew. Chem., Int. Ed. Engl. **1991**, 30, 1160. (b) More recently, see: Cremer, D.;

⁽⁶⁴⁾ A reviewer has suggested a plausible explanation for the S.: S bond shortening in solution. The lowest unoccupied orbital of the $[H_2S.:SH_2]^+$ complex (the S.: S σ^* orbital has one electron) involves a combination of the four S-H σ^* bonds with the right symmetry to overlap in a π fashion between the two sulfur atoms. The four explicit water molecules transfer charge into the LUMO which increases the S-H bond distances and decreases the S.: S distance.



Figure 4. van't Hoff plot for reaction 1.

Table 3. Calculated Entropies (J/K mol) and Enthalpies (kJ/mol) of Association at Various Temperatures^a

	298 K			520	520 K ^{b,c}		576 K ^d		
	ΔS	ΔH/a	Δ <i>H/</i> b	Δ <i>H</i> /c	ΔS	Δ <i>H</i> /a	ΔS	ΔH/a	Δ <i>H</i> /b
(SH ₂) ₂ ⁺	-120 ^e	-119	-128	-124					
$(DMS)_2^+$	-142^{f}	-118	-126				-133	-114	-122
$(DES)_2^+$	-151	-116			-143	-112			

^{*a*} Entropies are calculated at UHF/6-31G(d). Level a is PMP2/6-31G(d)//UHF/6-31G(d)+(0.9)ZPC/C_p//UHF/6-31G(d). Level b is [PMP4/6-31+G(2df,p)]//UMP2/6-31G(d)+(0.9)ZPC/C_p//UHF/6-31G(d). Level c is G2+C_p//UHF/6-31G(d). The symbol "(0.9)ZPC/C_p" indicates that the zero-point correction is made with a scaling factor of 0.9 and the heat capacity correction is made with a scaling factor of 0.9 and the reference monomers for (DES)₂⁺ are **3a**,**a**⁺. If **3d**,**d**⁺ are taken as the reference, the entropy and enthalpy of association become -150 J/K mol and -123 kJ/mol, respectively. ^c At 520 K the experimental entropy and enthalpy of association of (DES)₂⁺ are -132 J/K mol and -116 kJ/mol, respectively. ^d At 576 K the experimental entropy and enthalpy of association of (DMS)₂⁺ are -112 J/K mol and -111 kJ/mol, respectively. ^e The value of ΔS for (DMS)₂⁺ calculated at MP2/6-31G(d) is -120 J/K mol.

ion residence times increase with decreasing extraction voltage. As the Me_2S concentration is increased, the voltage dependence lessens and a limit is reached. Data as a function of total ion source pressure and concentration were extrapolated by linear least squares to zero extraction voltage. The extrapolated values from all the data lead to the van't Hoff plot in Figure 4.

The data from Figure 4 result in a reaction enthalpy of -111 \pm 2 kJ/mol (reported errors are one standard deviation) at the midtemperature of 576 K. Correcting the experimental data to 298 K using the computed geometry and vibrational frequencies results in an experimental association enthalpy of -115.0 kJ/ mol at that temperature. The corrected experimental association enthalpy (Table 4) agrees well with the ab initio-predicted association enthalpy of -126 kJ/mol (Table 3). The computed S-S bond distance of 2.795 Å can be compared to an average S-S single-bond distance of about 2.05 Å.⁶⁵ The bond energy (about 40% of a normal 2c-2e S-S bond⁶⁶), bond distance (about 135% of a normal 2c-2e S-S bond), and symmetric spin distribution between the two sulfur atoms all support the 2c-3e bonding picture. The geometry calculated at the MP2/6-31G-(d) level (Figure 2) is in good agreement with the geometry reported earlier at the UHF/3-21G(d) level.33

 Table 4.
 Experimental Heats of Formation and Association

 Enthalpies (kJ/mol) of 2c-3e Complexes^a

	heats of form	nation (kJ/mol)	association enthalpy (kJ/mol)		
	$\Delta H^{\circ}_{\mathrm{f},0}$	$\Delta H^{\circ}_{f,298}$	$\Delta H^{\circ}_{\rm rxn,0}$	$\Delta H^{\circ}_{\rm rxn,298}$	
$\frac{\overline{(\mathrm{SH}_2)_2^{+ b}}}{(\mathrm{DMS})_2^+}$	848.1 657.1	852.5 656.5	-119.6 -115.6	-123.7 -115.0	
$(DES)_2^+$	533.4	532.3	-120.4	-119.3	

^{*a*} Based on experimental heats of formation of neutral and cation monomer,³⁷ experimental enthalpies of association, and theoretical heat capacity corrections. ^{*b*} G2 enthalpy of association is used.

The geometry of $[Me_2S:SMe_2]^+$ (Figure 2) has C_{2h} symmetry. The tilt angle (the angle between the S-S bond and the line bisecting the C-S-C angle) is 95.6°. This is very similar to the calculated tilt angle in $[H_2S:SH_2]^+$ (92.3°). The S:S distance in $[Me_2S:SMe_2]^+$ is shorter than in $[H_2S:SH_2]^+$, possibly due to reduced electrostatic repulsion in $[Me_2S:SMe_2]^+$ since there are more atoms to delocalize the positive charge.

The intercept of the van't Hoff plot (Figure 4) results in $\Delta S^{\circ}_{rxn,576} = -112 \pm 3$ J/mol K, while the predicted entropy of association is $\Delta S^{\circ}_{rxn,576} = -133$ J/mol K. These results are in excellent accord with that expected for an ion-molecule association reaction resulting in a strongly bound cluster.⁶⁷ The negative entropy change is mostly due to the loss of three translational degrees of freedom upon reaction. The experimental value for the entropy of reaction is just outside the range which was assumed (-84 to -105 J/mol K) in our previous estimate for this bonding interaction. The computational results for the [Me₂S::SMe₂]⁺ system are summarized in Table 3 at 298 K and 576 K (the midexperimental temperature) and in Table 4 where the experimental thermodynamic properties have been corrected to 0 and 298 K. The present experimental free energy of reaction at 525 K (-52.2 kJ/mol) differs somewhat from a previous determination at the same temperature (-56.1)kJ/mol) using a different ion source.³³ This ion source probably suffered from poor thermalization of all the ion source parts, especially the ion source exit slit which was probably colder than the ion source body, consistent with the direction of the difference between the two measurements.

[Et₂S∴SEt₂]⁺. At low ionizing electron energies (near the threshold for formation of ions), three ions always present in the mass spectrum are Et₂S⁺, [Et₂S∴SEt₂]⁺, and [Et₂S∴SEt]⁺, where the latter ion was less intense than the monomer and dimer ions but still present under all explored experimental conditions. However, we were unable to definatively determine if [Et₂S∴SEt]⁺ originates from decomposition of the activated dimer complex or from electron impact fragmentation. Nevertheless we were able to establish equilibrium over a large temperature range (482–572 K). As the ionizing electron energy was increased, peaks which can be attributed to [Et₂S∴SCH₂Et]⁺ and [Et₂S₂]⁺ appeared in the mass spectrum. The ion intensity at *m*/z 165 due to [C₇H₁S₂]⁺ results from a clustering reaction of *m*/z 75 [C₃H₇S]⁺ with Et₂S.

The experimental equilibrium measurements on the $[Et_2S:.SEt_2]^+$ cluster were more challenging than those for $[Me_2S:.SMe_2]^+$. Under most conditions, RTD's showed that reaction 2 was close to equilibrium and started to merge as the extraction field was reduced (longer time). Figure 5 presents four distributions, two each for Et_2S^+ and $[Et_2S:.SEt_2]^+$ at different ion extraction fields. The figure clearly shows that the residence time distributions at shorter times are not identical but that the pair at the longer times are more similar. It was

^{(65) (}a) For references to S-S bond lengths, see: Knop, O.; Boyd, R. J.; Choi, S. C. J. Am. Chem. Soc. **1988**, 110, 7299. (b) Allen, F. H.; Kennard, O.; Watson, D. G.; Brammer, L.; Orpen, A. G.; Taylor, R. J. Chem. Soc., Perkin Trans. 2 **1987**, S1.

⁽⁶⁶⁾ Chang, R. Chemistry, 4th ed.; McGraw-Hill, Inc.: New York, 1991; p 378.

⁽⁶⁷⁾ Keesee, R. G.; Castleman, A. W., Jr. J. Phys. Chem. Ref. Data 1986, 15, 1011.



Figure 5. Residence time distributions for Et_2S^+ (open diamonds) and $[Et_2S \therefore SEt_2]^+$ (filled circles) at different ion extraction field strengths.



Figure 6. van't Hoff plot for reaction 2.

particularly difficult to reach equilibrium conditions at the lower temperatures. Much of our early data on reaction 2 did not fall on a linear van't Hoff line; rather it appeared that the data defined a limit for the line. As we gained experience with the chemical system, the van't Hoff plot in Figure 6 was collected. The data was obtained by extrapolating measured apparent equilibrium constants obtained as a function of ion extraction voltages to 0 V. Experiments at various total ion source pressures and composition were carried out. The two data points at the lowest temperatures in the van't Hoff plot shown in Figure 6 do not appear to fall on the line and were not included in the determination of the enthalpy and entropy of reaction. The least squares fit results in $\Delta H^{\circ}_{rxn,520} = -116 \pm 2$ kJ/mol and $\Delta S^{\circ}_{rxn,520} = -132 \pm 5$ J/mol K.

The monomers, SEt₂ and SEt₂⁺, are much more flexible molecules than SMe₂ and SMe₂⁺ and have multiple conformational minima. Three different structures were considered on the SEt₂ and SEt₂⁺ potential energy surface. The lowest energy structures (**3a**,**a**⁺) have C_{2v} symmetry and are in a trans-trans orientation with respect to rotation about the two S-C bonds (Figure 2). Two gauche-gauche structures were located; one pair were minima (**3b**,**b**⁺) while the other pair were transition states (**3c**,**c**⁺) at the HF/6-31G(d) level. Very close in energy to C_{2v} structures (**3a**,**a**⁺) are the gauche-trans structures of C_1 symmetry (**3d**,**d**⁺). At the PMP2/6-31G(d)+ZPC level, **3d** is only 0.2 kJ/mol higher in energy than **3a**, while **3d**⁺ is 1.4 kJ/ mol higher than **3a**⁺. Entropy favors the lower symmetries of **3d**,**d**⁺ over **3a**,**a**⁺. As shown in Table 5, the free energy of **3d**

Table 5. Calculated Relative Free Energies (kJ/mol) of DES and DES⁺ Monomers^{*a*}

	29	8 K	52	0 K
sym	DES	DES ⁺	DES	DES ⁺
C_{2v} 3a,a ⁺	0.0	0.0	0.0	0.0
C_2 3b,b ⁺	-0.1	5.4	-0.2	8.0
C_1 3d,d ⁺	-1.5	0.6	-2.7	1.2

^a $\Delta G(\text{rel to } C_{2v}) = [\Delta E(\text{PMP2}) + (0.9)\text{ZPC} + C_p(\text{corr})] - T\Delta S.$

is more favorable than that of **3a** by 1.5 kJ/mol at 298 K and by 2.7 kJ/mol at 520 K. The free energy of **3a**⁺ is slightly lower than that of **3d**⁺ but only by 1.2 kJ/mol at 520 K. Therefore, there exists the question about the appropriate reference monomers in calculating the binding energies. The free energy at 520 K favors the SEt₂ conformer **3d**, while it slightly favors the SEt₂⁺ conformer **3a**⁺. We have chosen to take **3a**,**a**⁺ as the reference since they are the lowest energy conformers at 0 K.

The conformation surface of a related neutral molecule, divinyl sulfide, has been previously studied.⁶⁸ The two stable conformations, identified from a study of the potential energy surface at the HF/3-21G(d) level, could both be described as gauche—gauche conformations.^{68a}

Two alternative 2c-3e complexes were considered for $[Et_2S:SEt_2]^+$. The complexes most analogous to $[H_2S:SH_2]^+$ and $[Me_2S::SMe_2]^+$ is the $[Et_2S::SEt_2]^+$ complex with C_{2h} symmetry. The $[Et_2S::SEt_2]^+$ complex can be formed from $[Me_2S::SMe_2]^+$ by replacing the hydrogen on Me_2S which points away from the opposite Me₂S with a methyl group. This replacement should give a structure which minimizes intermoiety steric repulsion in the complex between Et_2S , 3c, and SEt_2^+ , $3c^+$ but which may not minimize *intramoiety* repulsion (within each monomer unit in the $[Et_2S::SEt_2]^+$ complex). This is particular evident as the $3c,c^+$ structures are 5.0 and 7.1 kJ/ mol higher in energy than $3a_{,a}^{+}$. A complex formed between $3a_{,a}^{+}$ (not calculated) would have more intermoiety but less intramoiety repulsions. A compromise complex formed between 3d,d⁺ would have the advantage of having a very modest "cost" to adopt the higher energy conformer (3d,d⁺ is 1.7 kJ/mol higher in energy than $3a_{,a^+}$) yet not have unfavorable intermoiety repulsion. The complex between $3d_{,d}^{+}$ has C_2 symmetry and is 7.0 kJ/mol lower in energy than the C_{2h} symmetry complex. Across the S: S 2c-3e bond, two CH_2 groups are trans with respect to each other while the other two CH₂ groups are cis. The methyl groups are gauche-trans about the S-C bonds within each moiety, just as they are in 3d,d⁺.

The ab initio calculations at the PMP2/6-31G(d)//HF/6-31G(d) level predict a bond energy of 117.0 kJ/mol. At 298 K, ΔH°_{TXN} and ΔS°_{TXN} are -116 kJ/mol and -151 J/mol K, respectively. Correcting to 520 K, the thermodynamic values become -112 kJ/mol and -143 J/mol K which can be compared to the experimental values of -116 kJ/mol and -132 J/mol K, respectively. The predicted [Et₂S.:SEt₂]⁺ geometry (Figure 2) has a S-S bond distance of 2.870 Å. Using the calculated molecular parameters to correct the experimental enthalpies to 0 K results in an experimental bond energy of 120.4 kJ/mol for [Et₂S.:SEt₂]⁺ (Table 4).

Comparison of Calculated Properties of $[H_2S \therefore SH_2]^+$, $[Me_2S \therefore SMe_2]^+$, and $[Et_2S \therefore SEt_2]^+$. 1. Comparison of Monomer Ionization Energies. A comparison of adiabatic ionization energies is made with experiment in Table 6. As the level of theory improves, the difference between theory and experiment

^{(68) (}a) Kimmelma, R.; Hotokka, M. *THEOCHEM* **1993**, 285, 71. (b) Kao, J.; Eyermann, C.; Southwick, E.; Leister, D. J. Am. Chem. Soc. **1985**, 107, 5323.

Table 6. Calculated Adiabatic Ionization Energies (eV)^{a,b}

	IE/a	IE/b	IE/c
SH₂ DMS DES	9.82 (0.63) 8.24 (0.45) 7.99 (0.44)	10.18 (0.27) 8.52 (0.17)	10.43 (0.02)

^a Level a is PMP2/6-31G(d)//UHF/6-31G(d)+(0.9)ZPC. Level b is [PMP4/6-31+G(2df/p)]//UMP2/6-31G(d)+(0.9)ZPC. Level c is G2. ^b The values in parentheses are the difference between theory and experiment. The experimental values for SH₂, DMS, and DES are 10.45, 8.69, and 8.43 eV, respectively. Lias, S. G.; Bartmess, J. E.; Liebman, J. F.; Holmes, J. L.; Levin, R. D.; Mallard, W. G. J. Phys. Chem. Ref. Data **1988**, 17 (Suppl. 1). A similar value for the IE of DMS (8.64 eV) has recently been reported; see: Nourbakhsh, S.; Norwood, K.; Yin, H.-M.; Liao, C.-L.; Ng, C. Y. J. Chem. Phys. **1991**, 95, 5014.

Table 7. Calculated Vibrational Frequencies (cm^{-1}) at the UHF/6-31G(d) Level Due to the Association of Two Monomers^{*a*}

mode	$(SH_2)_2^+ (C_{2h})$	$(DMS)_2^+ (C_{2h})$	$(DES)_2^+(C_2)$
S∴S stre	232 (230) ag	283 (274) ag	202 a
S∴S rock	587 (527) a _g	194 (187) a _g	182 a
S∴S rock	441 (427) b _u	194 (188) b _u	184 Ь
S∴S rock	456 (416) a _u	176 (161) a _u	98 a
S∴S rock	533 (470) b _g	175 (164) b _g	91 b
S∴S tors	66 (75) a _u	29 (32) a _u	19 a

^a Values given in parentheses are calculated at MP2/6-31G(d).

Table 8. Calculated λ_{max} (transitions in nm) for Various 2c-3e Bonded Species^{*a*}

	//UHF/6-31G(d)	//UMP2/6-31G(d)	exptl ^b
$(SH_2)_2^{+}4H_2O$	368	351	370
$(SH_2)_2^+$	405	396	370
$(DMS)_2^+$	476	452	465
$(DES)_2^+$	509		485

^{*a*} Excitation energies were calculated with single excitations (CIS/ 6-31G(d)). ^{*b*} The experimental λ_{max} are measured in aqueous solution.²⁵

decreases. The difference in the ionization energy of SH_2 between G2 and experiment is only 0.02 eV.

2. Comparison of New Frequencies. The new frequencies associated with the relative motions of the two sulfides in the 2c-3e complex are reported at the UHF/6-31G(d) level in Table 7 (MP2/6-31G(d) for $[H_2S::SH_2]^+$ and $[Me_2S::SMe_2]^+$). As one proceeds from $[H_2S::SH_2]^+$ to $[Me_2S::SMe_2]^+$ to $[Et_2S::SEt_2]^+$, one finds the frequencies of associations (HF/ 6-31G(d)) to become smaller which simply reflects the increased reduced mass of the sulfide. The S: S stretch remains relatively constant at about 200-280 cm⁻¹. In contrast, the torsional mode decreases from 66 to 29 to 19 cm⁻¹ as one proceeds from $[H_2S \therefore SH_2]^+$ to $[Me_2S \therefore SMe_2]^+$ to $[Et_2S \therefore SEt_2]^+$. Since this mode is present in the product but not in reactants, entropy and heat capacity calculations will be very dependent on how accurately this mode is calculated. For example, an overestimation of the 19 cm^{-1} mode by only 10 cm^{-1} , would reduce the contribution by approximately 6.2 J/mol K at 298 K.

3. Comparison of λ_{max} . Since the magnitude of the splitting between the σ and σ^* orbitals is related to the degree of interaction between the two moieties forming the 2c-3e bond, λ_{max} is sometimes taken as an indicator of 2c-3e bond strength. On the basis of a series of measurements in aqueous solution, Asmus found^{9-21,25} that the λ_{max} varied in a manner which could be related to the Taft σ^* value.²⁵ The values of λ_{max} for $[H_2S \therefore SH_2]^+$, $[Me_2S \therefore SMe_2]^+$, and $[Et_2S \therefore SEt_2]^+$ increase in the order 370, 465, and 485 nm which would indicate decreasing 2c-3e bond strength in that order (Table 8).

However, Clark³⁰ and Bally²⁶ have pointed out that alkyl substituents produce a red shift "independent of the steric interactions". The λ_{max} transition occurs from an orbital that

has strong σ lone pair (perpendicular to S: S axis) character to an orbital that has mostly p lone pair (parallel to S: S axis) character. Alkyl substituents have the right symmetry to destabilize the σ lone pair orbital more than the p lone pair orbital resulting in a smaller energy gap and a red shift.

In order to explore the relationship between the optical transition and S..S bond energies, we carried out theoretical calculations of the low-lying transitions. An electron correlation method of including single excitations was used⁵⁸ with the 6-31G(d) basis set (CIS/6-31G(d)) to determine the lowest three states. In each case, the lowest transition corresponded to an excitation from the doubly occupied lone pair combination orbital to the singly occupied sulfur-sulfur σ^* orbital.²⁶ Thus, the excitation was a ${}^{2}B_{u} \rightarrow {}^{2}Ag$ except for $[Et_{2}S..SEt_{2}]^{+}$ which was a ${}^{2}B \rightarrow {}^{2}A$. The experimental trend of the first transition energy is reproduced (red shift) even though the S..S bond strength remained essentially constant, varying less than 10 kJ/ mol for the series.

Since λ_{max} was measured in aqueous solution, we were also interested in estimating the effect of solvation on λ_{max} , especially given the fact that the S::S distance is predicted to decrease upon solvation (see above). In order to investigate this possibility, we calculated the lowest excitation energy for the [(SH₂)₂]⁺•4H₂O complex. As discussed above, the four waters shorten the S::S bond which results in a decrease of the wavelength needed for excitation (larger transition energy) by 37 nm at the CIS/6-31G(d)//UHF/6-31G(d) level and 45 nm at the CIS/6-31G(d)//MP2/6-31G(d) level. We conclude that solvation effects should cause a λ_{max} blue shift and that substituents which destabilize the σ lone pair in the 2c-3e complex should cause a red shift. We caution that extrapolation of λ_{\max} trends to S: S bond strength trends should be made with care since a blue shift (solvent induced) and a red shift (substituent induced) can occur without significantly affecting the S:S bond strength.

4. Comparison of $a_{\rm H}$ Hyperfine Coupling Constants. Another point of contact between experiment and theory for 2c-3e bonded S: S complexes is through $a_{\rm H}$ hyperfine coupling constants which have been reported for hydrogen in $[{\rm Me}_2{\rm S}: {\rm SMe}_2]^+$ and $[{\rm Et}_2{\rm S}: {\rm SEt}_2]^+$ complexes.⁶⁹⁻⁸⁶ While the majority of unpaired spin is located evenly between the two

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Table 9. Comparison of Calculated and Experimental $a_{\rm H}$ Hyperfine Coupling Constants (G)^{a,b}

=	-			
	UHF/6-31G(d)//a	UHF/6-31G(d)//b	UHF/6-31+G(2df,p)//b	exptl
SH ₂ ⁺ (2H)	-33.4	-33.4	-26.1	
SH ₂ +•2H ₂ O (2H)	-30.8	-30.1	-25.8	
$(SH_2)_2^+(4H)$	-25.1	-20.1	-17.1	
$(SH_2)_2^{+}4H_2O(4H)$	-19.5	-19.0	-16.4	
DMS ⁺ (6H)	16.7	18.0	16.9	20.4 (6H) ^c
$(DMS)_2^+$ (12H)	6.6	6.4	6.0	$6.6 (12H)^d$
DES ⁺ 3a ⁺ (4H)	22.3			18-20 (2H) ^e
DES ⁺ 3b ⁺ (2H)	10.1			18-20 (2H) ^e
DES ⁺ 3d ⁺ (4H)	15.1			18-20 (2H) ^e
$(DES)_2^+ C_{2h} (8H)$	3.0			6.9 (8H) ^c
$(DES)_2^+ C_2 (8H)$	6.9			6.9 (8H) ^d

^a Level a indicates geometry optimization at UHF/6-31G(d), and level b indicates geometry optimization at UMP2/6-31G(d). ^b The Fermi contact integral is averaged over the number of hydrogens indicated in parentheses and multiplied by 1600.^{87,88} ^c Reference 74. ^d Reference 83. ^e Reference 78.

sulfurs, sufficient spin density exists on hydrogens to give rise to a detectable ESR spectrum. From a computational point of view, $a_{\rm H}$ can be calculated by averaging the Fermi contact integral over the appropriate number of symmetry-equivalent hydrogens and multiplying by the constant 1600.87,88 These results are reported in Table 9 where the Fermi contact integral is calculated with the UHF spin density. It is known that much improved results can be obtained when the Fermi contact integral is calculated with the UMP2 spin density.^{87,88} However, our intention is only to determine whether the qualitative trends can be reproduced by theory. The answer to this question is yes. The changes in the value of $a_{\rm H}$ between the monomer and the 2c-3e bonded complex are well reproduced. The effect of the medium on the calculated $a_{\rm H}$ values can be determined by comparing calculations on $[H_2S:SH_2]^+$ with and without explicit waters. Using the 6-31+G(2df,p) basis set, the difference in calculated values of $a_{\rm H}$ is less than 1 G.

IV. Conclusions

Experimental gas-phase bond energies have been determined for the 2c-3e S \therefore S bond in [Me₂S \therefore SMe₂]⁺ and [Et₂S \therefore SEt₂]⁺. Chemical equilibrium was established over a temperature range of 533-626 K for reaction 1 and over a range of 482-572 K for reaction 2. From van't Hoff plots, a value of 111 kJ/mol was determined for the bond enthalpy of the S \therefore S bond in [Me₂S \therefore SMe₂]⁺ while a value of 116 kJ/mol was determined for the same bond in [Et₂S \therefore SEt₂]⁺. Ab initio calculations on the 2c-3e bonded complexes, [H₂S \therefore SH₂]⁺, [Me₂S \therefore SMe₂]⁺, and $[Et_2S:SEt_2]^+$, are in good agreement with experiment. At the highest level of theory, the S:S bond energies are 120 kJ/ mol (G2), 126 kJ/mol ([PMP4/6-31+G(2df,p)]), and 117 kJ/ mol (PMP2/6-31G(d)), respectively. These binding energies are in the middle range of ion-molecule cation association enthalpies which vary from less than 20 kJ/mol to more than 200 kJ/mol.⁶⁷

From a combination of experimental and theoretical data, we have estimated the heats of formation of the three complexes at 298 K to be 852.5, 656.5, and 532.3 kJ/mol, respectively. Other calculated properties include adiabatic ionization energies (for neutral monomers), excitation energies, and hydrogen hyperfine coupling constants. The wavelength value of the excitation energy agrees with the experimental trend observed by Asmus in solution.

Three conformers of Et_2S and Et_2S^+ were identified computationally. The most stable $[Et_2S::SEt_2]^+$ complex (C_2 symmetry) occurs from the association of C_1 conformers (gauche-trans) of Et_2S and Et_2S^+ . This complex has the best compromise of minimum intermoiety steric repulsion while maintaining a near optimal geometry of the monomers.

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